

Noise sensitivity for the 2D Stochastic Heat Equation and directed polymers

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University of Milano-Bicocca

Workshop on “Workshop on Irregular Stochastic Analysis ”

INdAM Meeting at Palazzone di Cortona ~ 24 June 2025

Collaborators



Joint work with
Anna Donadini

Collaborators



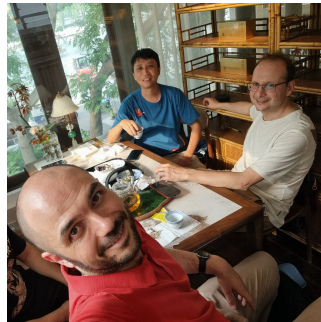
Joint work with

Anna Donadini

Previous works with

Rongfeng Sun (NUS)

Nikos Zygouras (Warwick)



Outline

1. The critical 2D SHF
2. Which equation for the SHF?
3. Noise sensitivity

The Stochastic Heat Equation

Heat equation with multiplicative singular potential

$$t \geq 0, x \in \mathbb{R}^d$$

$$\partial_t u(t, x) = \Delta_x u(t, x) + \beta u(t, x) \xi(t, x) \quad (\text{SHE})$$

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Natural candidate solution: the critical 2D Stochastic Heat Flow (SHF)

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Convergence of $u_N(t, \varphi) = \int_{\mathbb{R}^2} u_N(t, x) \varphi(x) dx$ as $N \rightarrow \infty$?

Renormalisation

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Then u_N converges in law to a **unique** and **non-trivial limit** \mathcal{U}^ϑ

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\mathcal{U}^ϑ = critical 2D **Stochastic Heat Flow (SHF)** = stochastic process of random measures on \mathbb{R}^2

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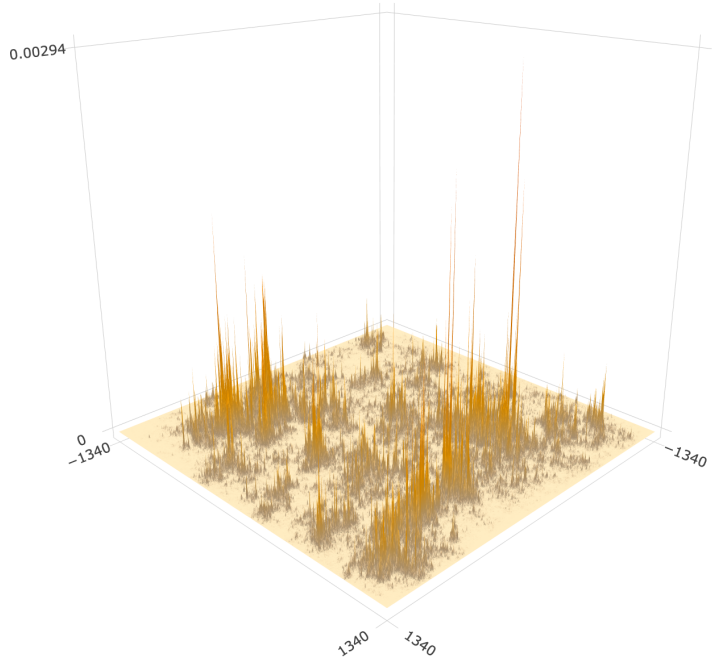
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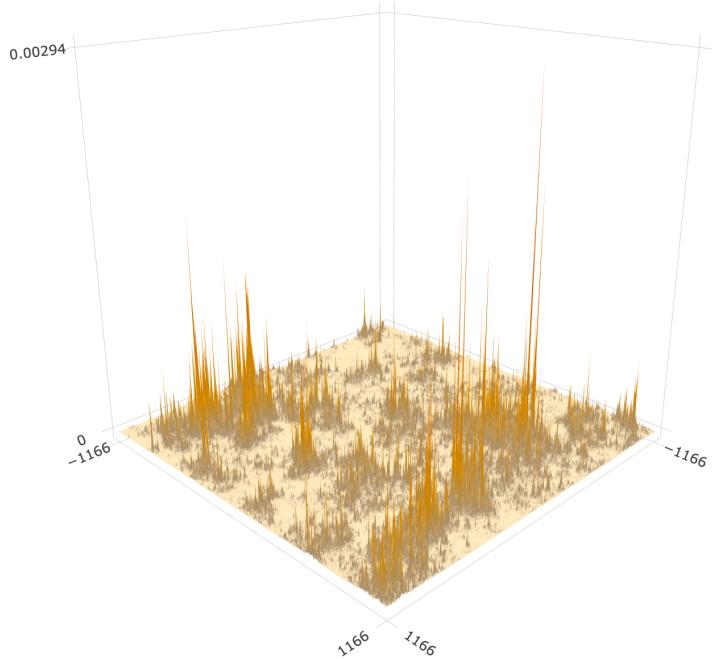
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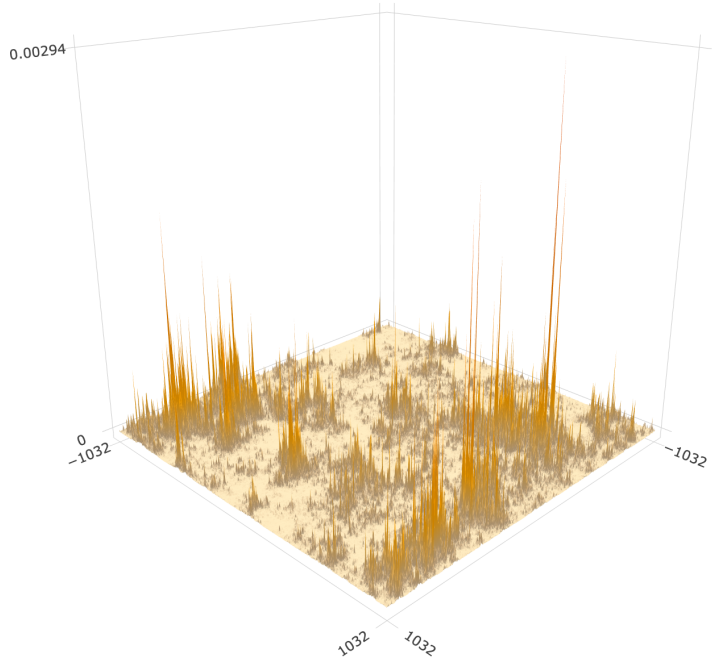
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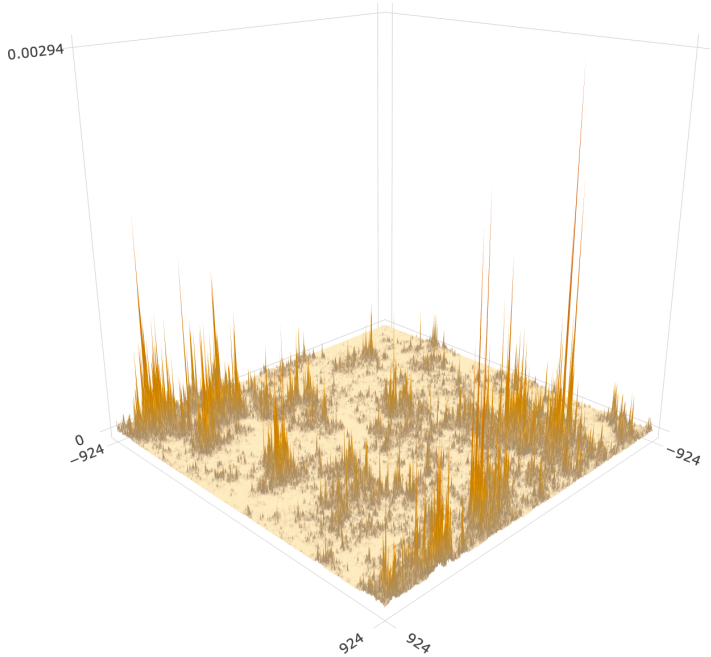
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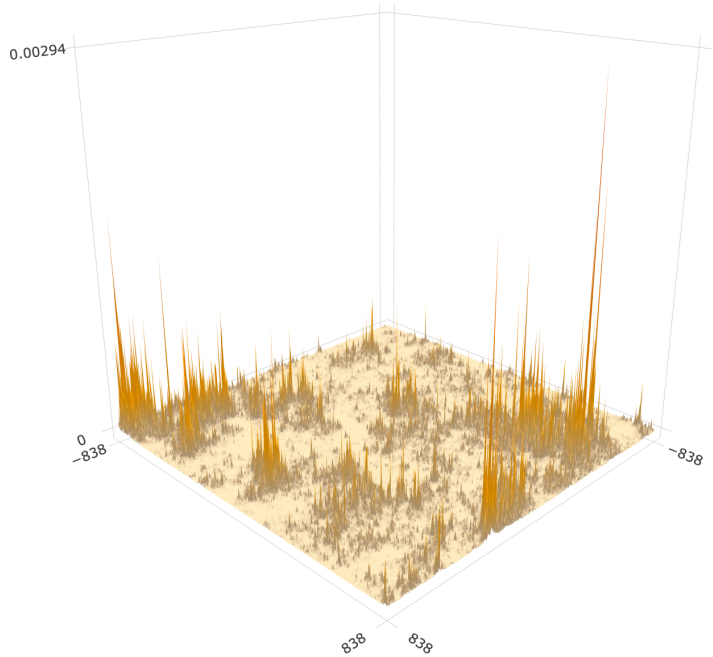
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- ▶ **Universality** w.r.t. approximation scheme [C.S.Z. 23] [Tsai 24+]

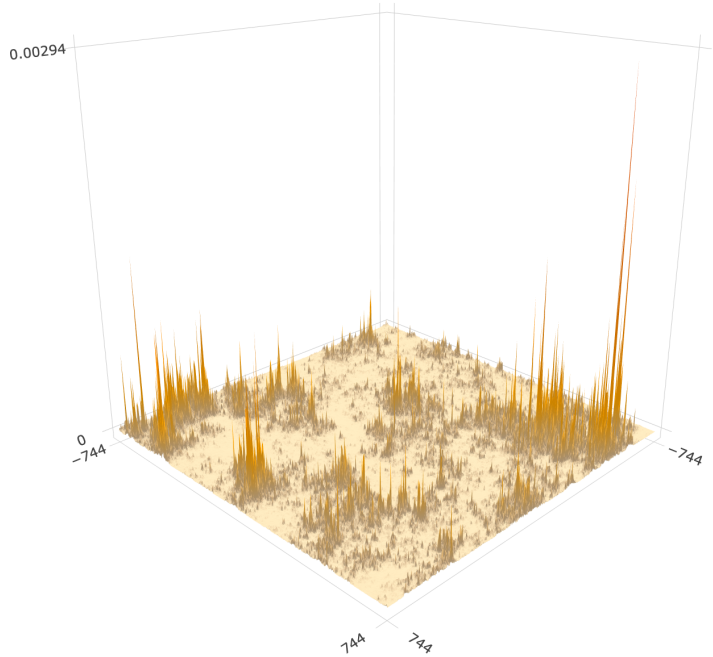


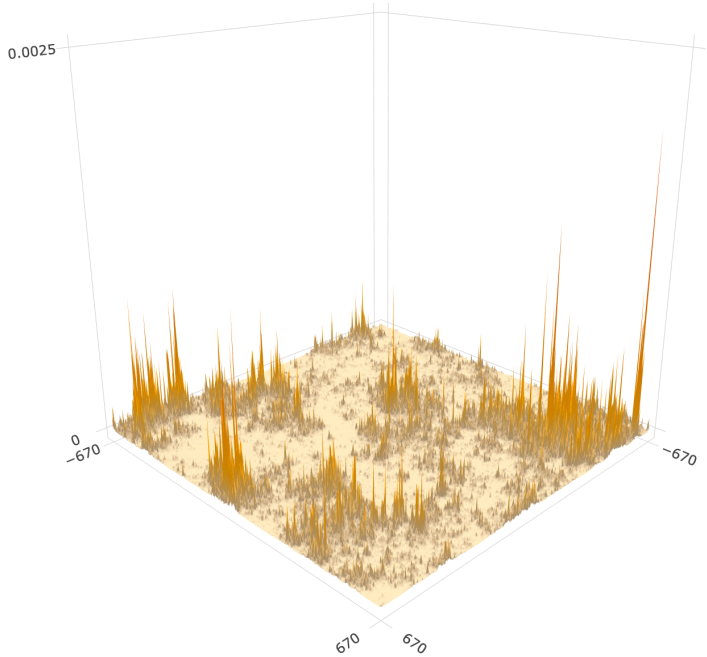


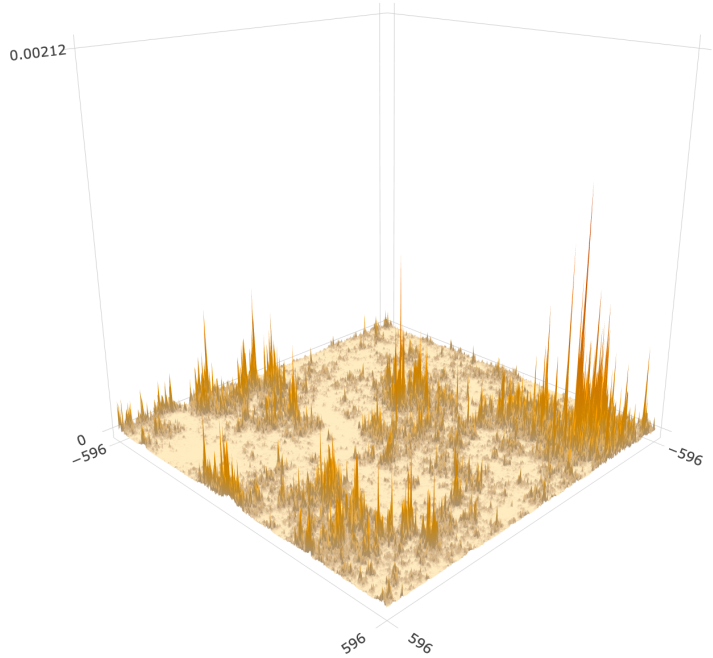


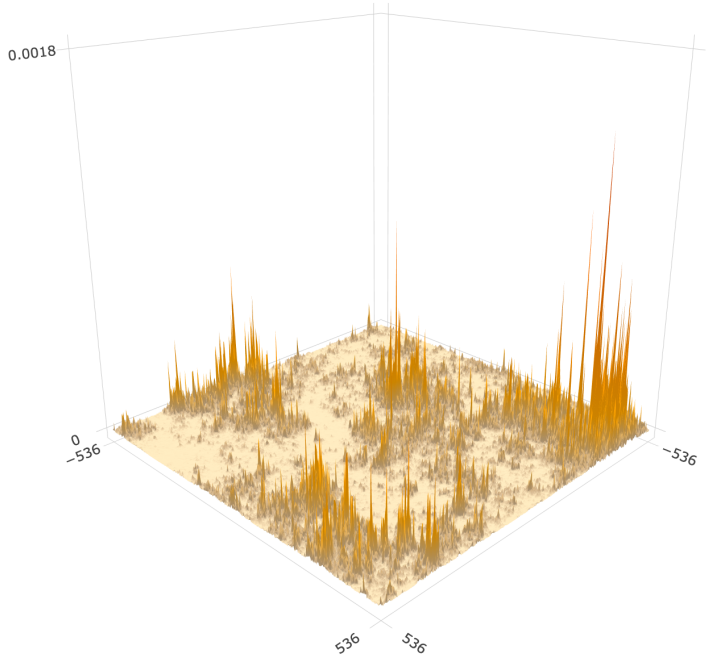


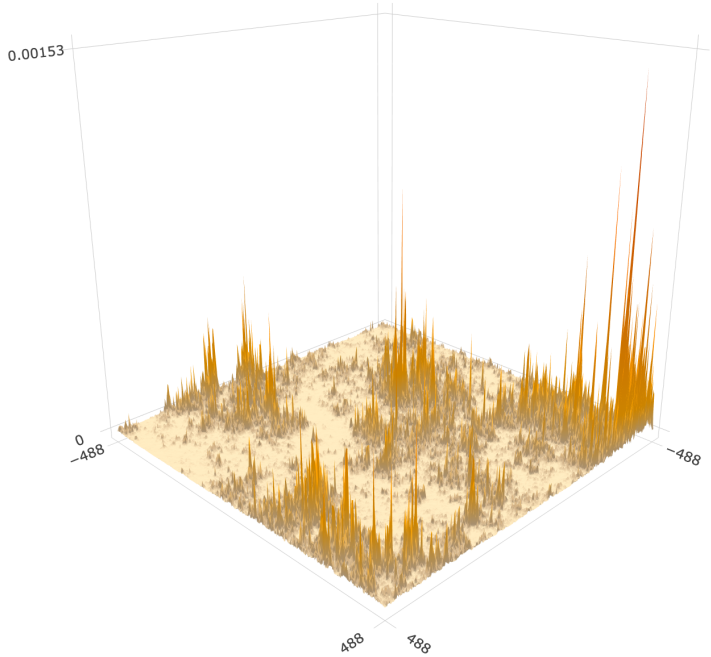


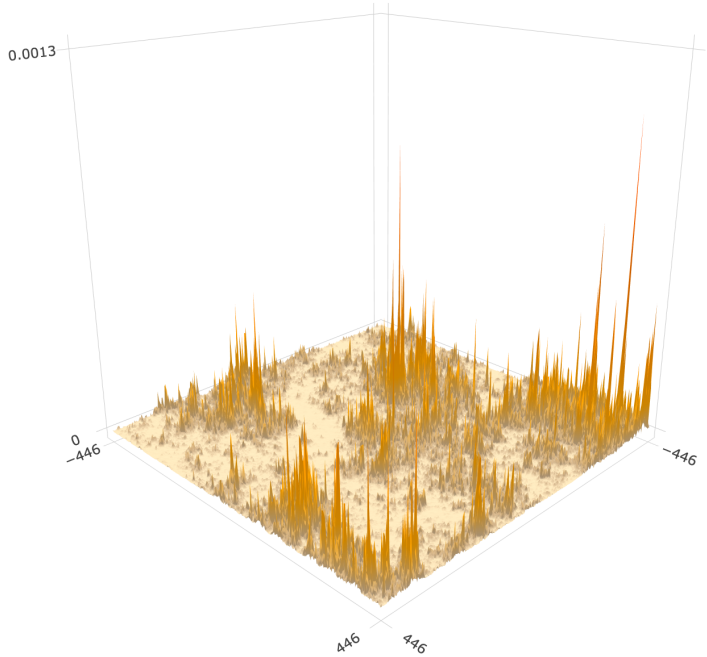


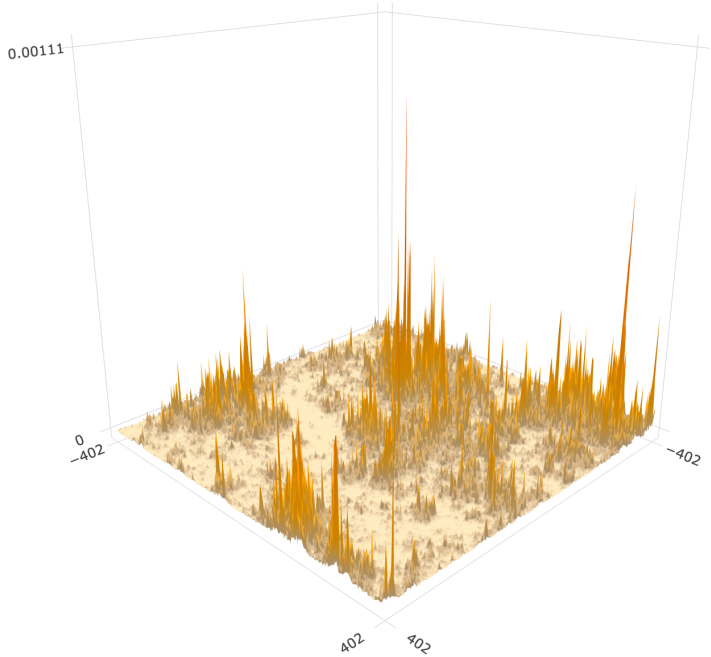


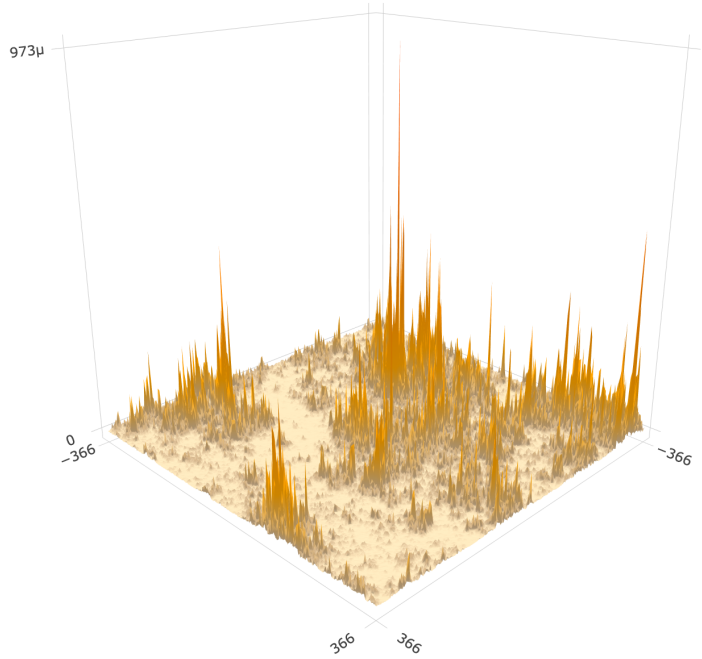


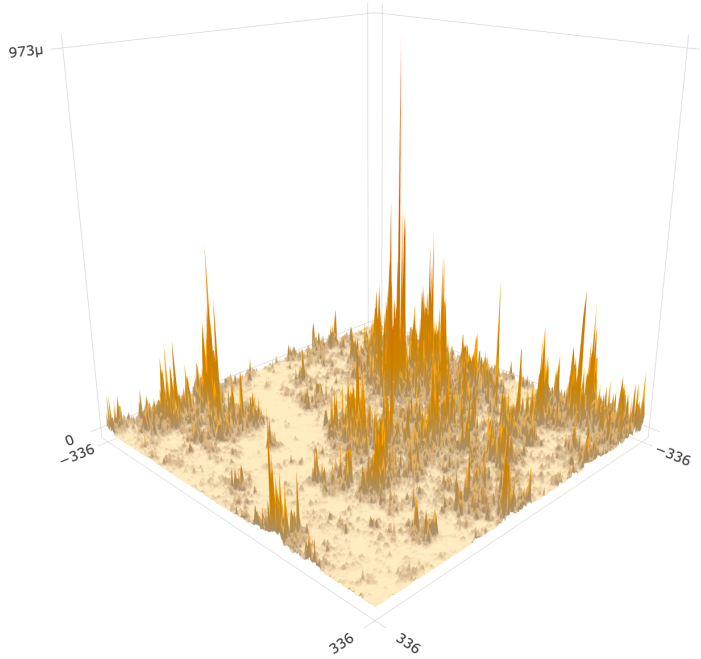


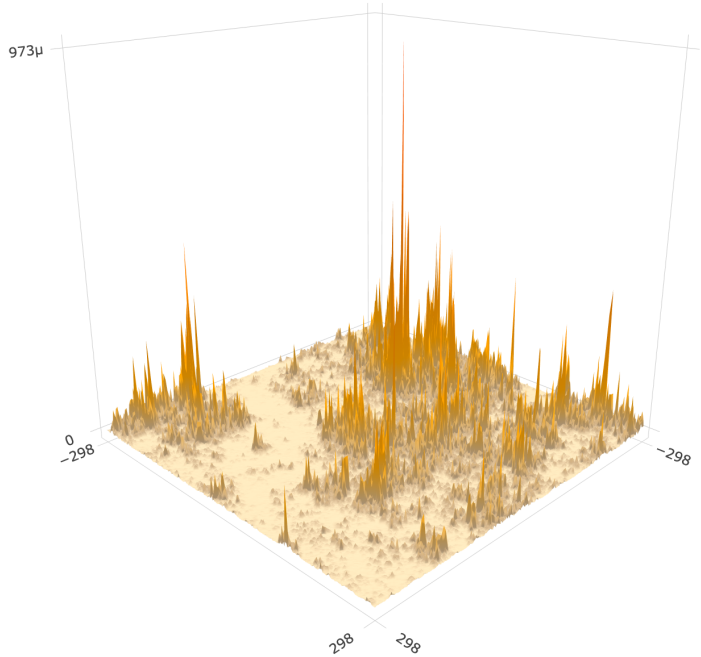


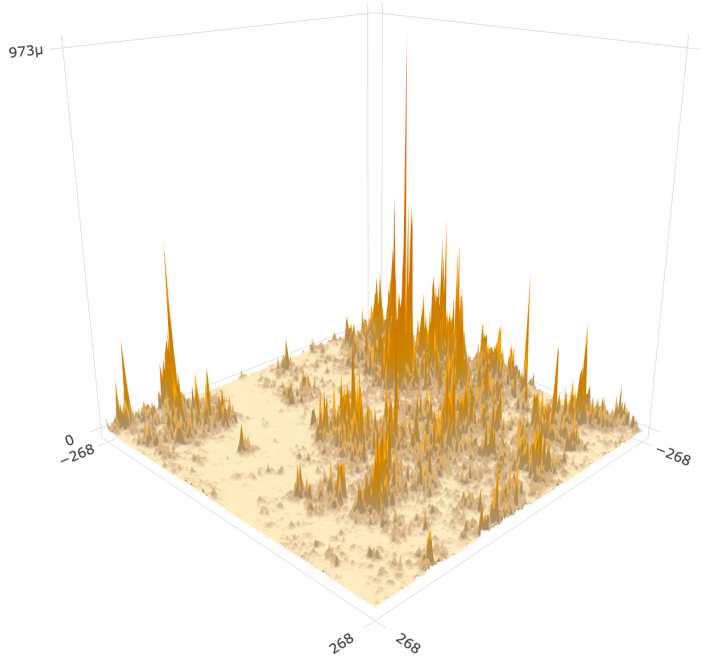


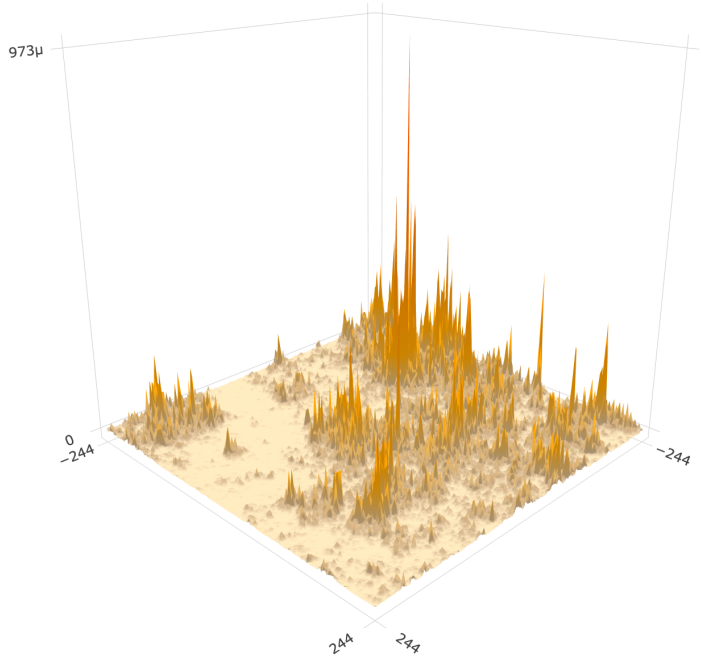


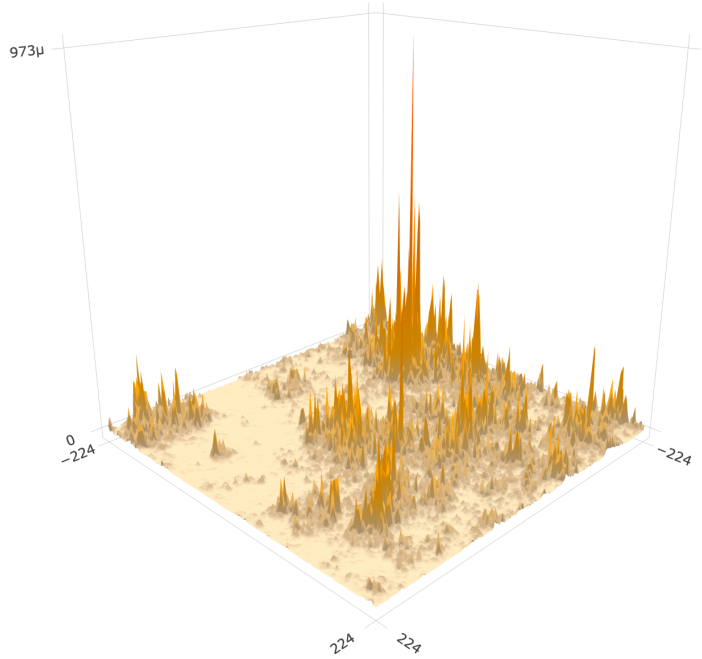


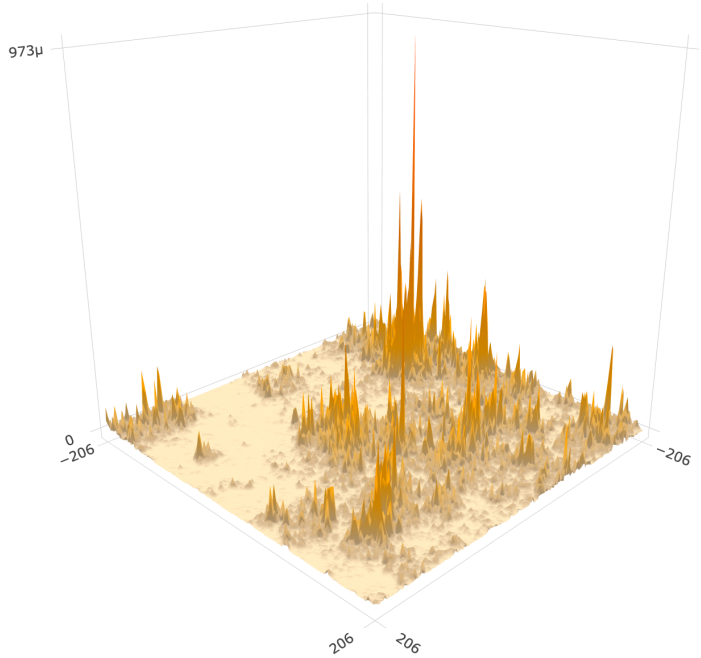


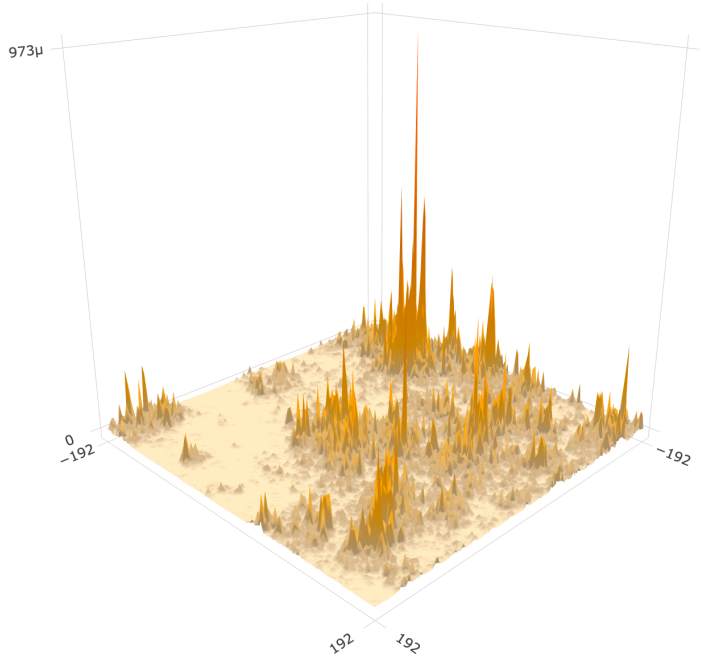


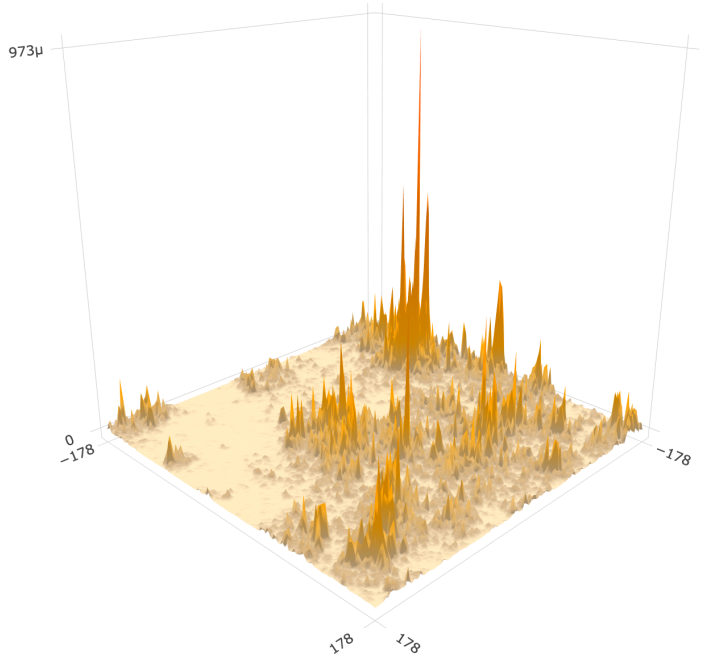


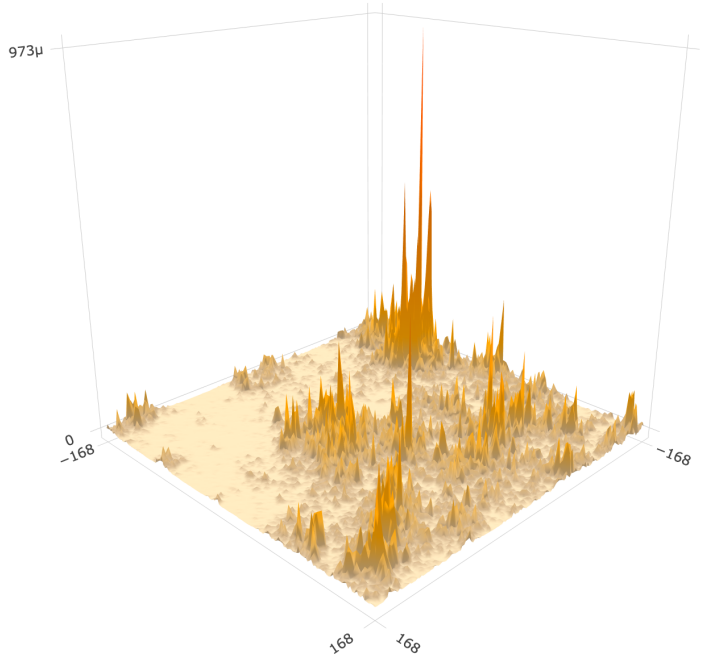


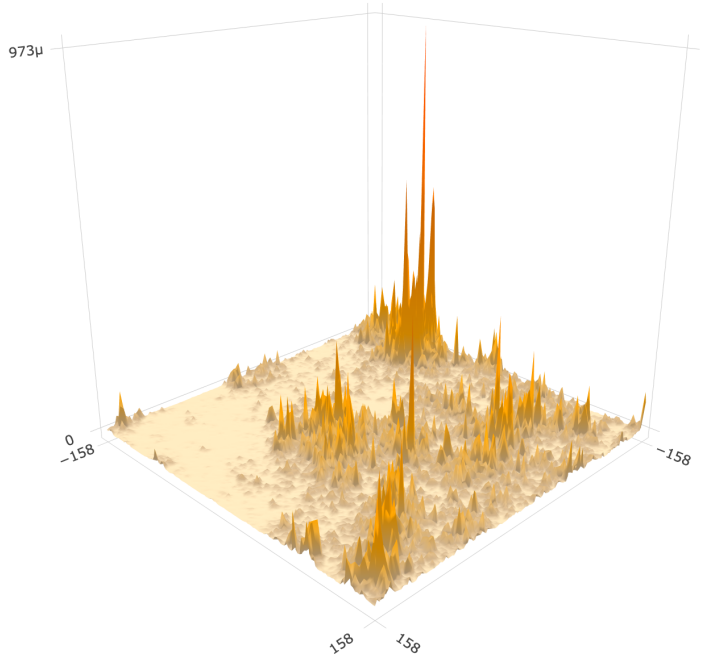


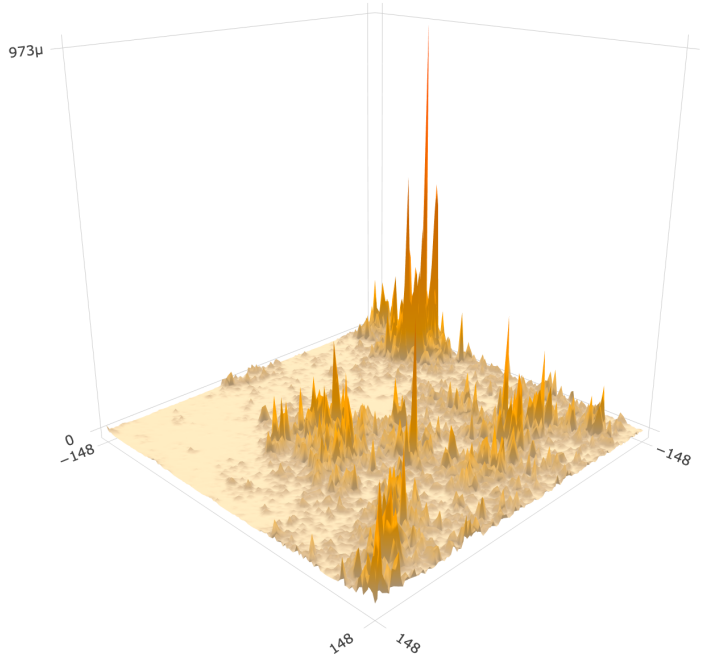


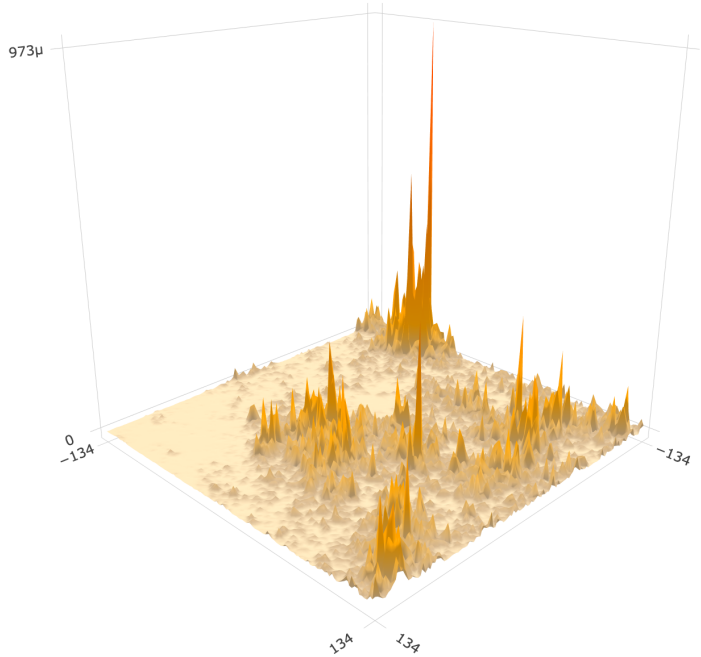


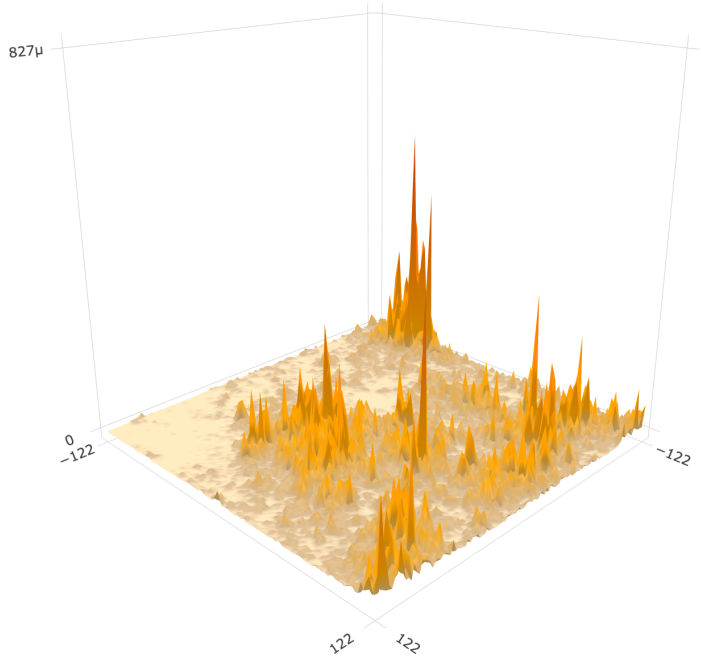


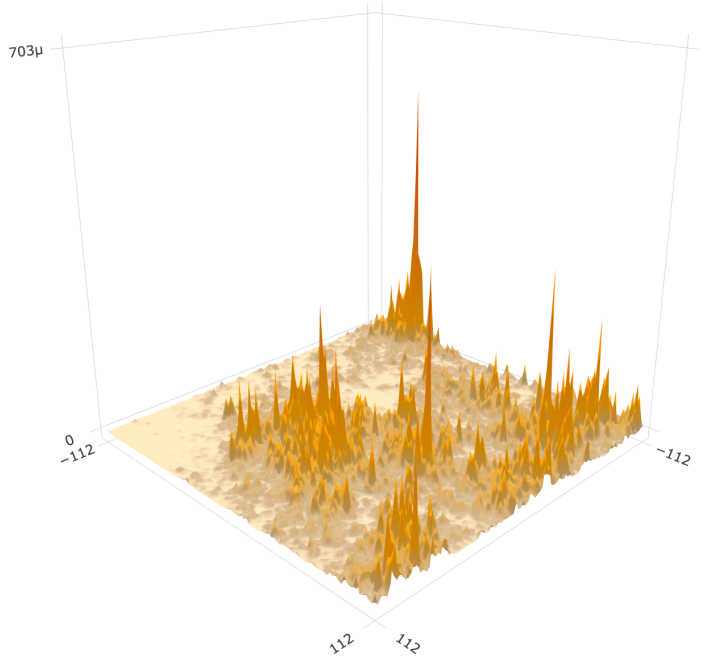


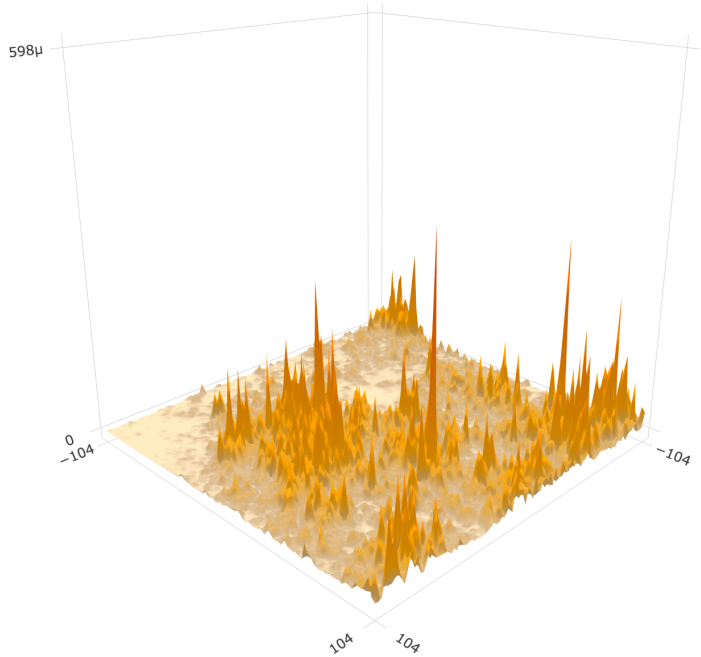


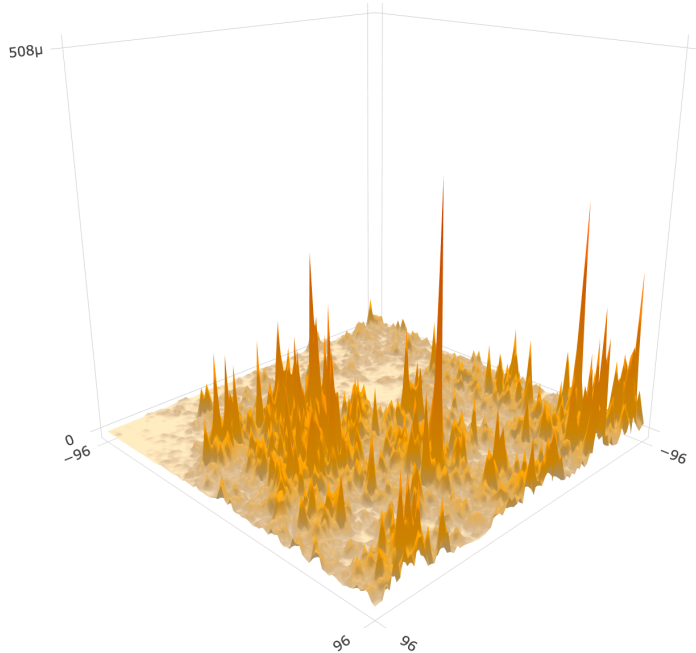


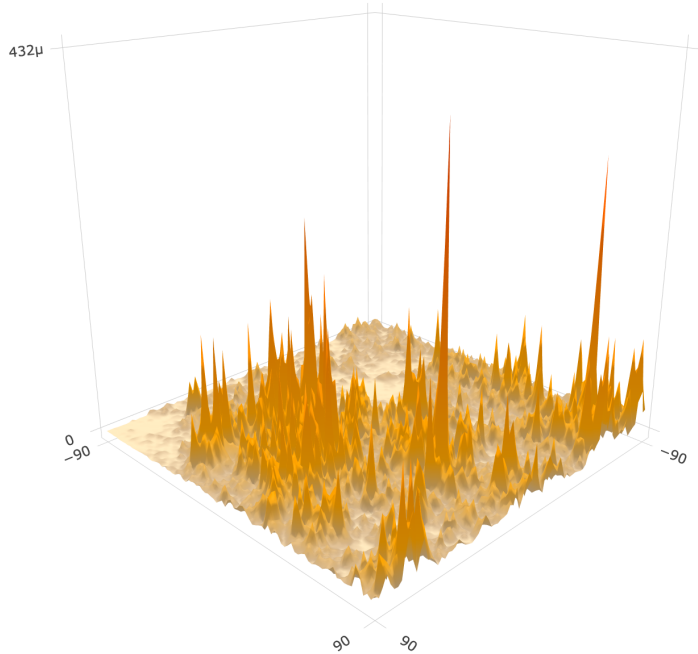












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SHF and the white noise

Does SHF $\mathcal{U}^\vartheta(t, dx)$ satisfy a SPDE driven by white noise $\xi(t, x)$?

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$$\langle \xi, \psi \rangle = \lim_{N \rightarrow \infty} \int \xi_N(t, x) \psi(t, x) dx \quad \text{in distribution} \quad \psi \in C_c^\infty(\mathbb{R}^{1+2})$$

No equation for the SHF

Theorem

[C.–Donadini 25+]

$$(\xi_N, u_N) \xrightarrow[N \rightarrow \infty]{d} (\xi, \mathcal{U}^{\vartheta})$$

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Recently proved: \mathcal{U}^ϑ is a “black noise” (à la Tsirelson) [Gu-Tsai arXiv 25]

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(partition function of 2D directed polymer in random environment)

Outline

1. The critical 2D SHF
2. Which equation for the SHF?
3. Noise sensitivity

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[Garban–Steif 14] [O’Donnell 14]

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$$\iff \phi(f_N(\omega^\varepsilon)) \text{ is noise sensitive}$$

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Spectral criterion

$$\text{Noise sensitivity} \iff \forall d \in \mathbb{N}: \|f_N^{(d)}\|_2^2 \xrightarrow{N \rightarrow \infty} 0$$

The BKS Theorem

Boolean setting: binary functions $f(\omega)$ of binary variables ω_i

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Robust condition for noise sensitivity based on influences

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$\forall \varepsilon > 0$: $\text{Cov}[f(\omega^\varepsilon), f(\omega)] \leq C \mathcal{W}(f)^{\alpha \varepsilon}$ [Keller–Kindler 13]

Influences beyond the Boolean setting

Define $\delta_i f := f - \mathbb{E}_i[f]$ with $\mathbb{E}_i[\cdot] = \mathbb{E}[\cdot | \sigma(\omega_j : j \neq i)]$ [Talagrand 94]

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It is the L^1 influence that is relevant for us: $\mathcal{W}(f) := \sum_i I_i^{(1)}(f)^2$

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Back to SHE

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$\implies u_N(t, \varphi)$ is asymptotically **independent** of any bounded order chaos

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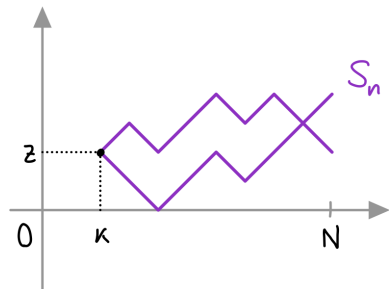
$$\mathbb{Cov} [f(\omega^\varepsilon), f(\omega)] \leq \mathcal{W}(f)^{\frac{\varepsilon}{2-\varepsilon}} + o(1)$$

The assumption that ω_i 's take finitely many values can hopefully be removed

Grazie

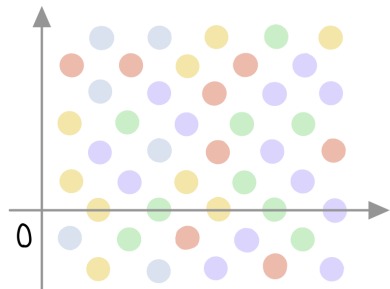
Directed Polymer in Random Environment

- ▶ $S = (S_n)_{n \geq 0}$ simple random walk on \mathbb{Z}^d



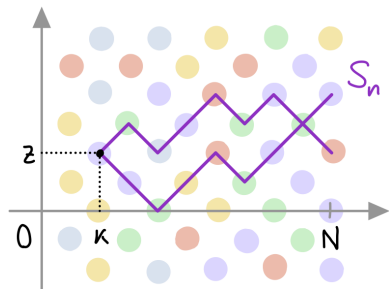
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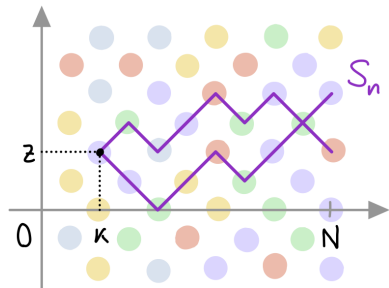
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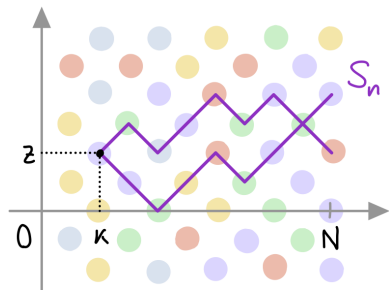
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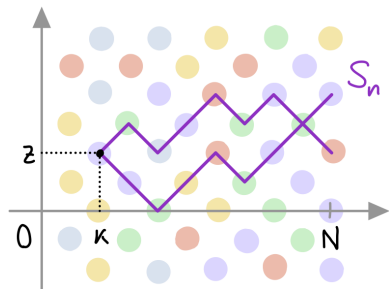
Partition Functions

$(k \in \mathbb{N}, z \in \mathbb{Z}^d)$

$$Z_{N,\beta}^{\omega}(k, z)$$

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Partition Functions

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$$Z_{N,\beta}^{\omega}(k, z) = \mathbb{E} \left[e^{\beta H(S, \omega) - \frac{1}{2} \beta^2 (N-k)} \mid S_k = z \right]$$

Partition functions and SHE

Diff. rescaled partition functions

$$Z_{N,\beta}^{\omega}(N(1-t), \sqrt{N}x) \quad (\text{time rev.})$$

Partition functions and SHE

Diff. rescaled partition functions = discretized SHE solutions

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Partition functions solve a difference equation:

with $\xi_N \approx \omega$

$$\begin{cases} \partial_t u_N(t, x) = \Delta_x u_N(t, x) + \underbrace{\beta N^{\frac{2-d}{4}}}_{\text{}} u_N(t, x) \xi_N(t, x) \\ u_N(0, x) \equiv 1 \end{cases} \quad (\text{reg-SHE})$$

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$$\begin{cases} \partial_t u_N(t, x) = \Delta_x u_N(t, x) + \underbrace{\beta N^{\frac{2-d}{4}}}_{\beta_{\text{SHE}}} u_N(t, x) \xi_N(t, x) \\ u_N(0, x) \equiv 1 \end{cases} \quad (\text{reg-SHE})$$

Partition functions and SHE

Diff. rescaled partition functions = discretized SHE solutions

$$Z_{N,\beta}^{\omega}(N(1-t), \sqrt{N}x) = u_N(t, x) \quad (\text{time rev.})$$

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Discrete analogue of Feynman-Kac

$$u_N(t, x) \approx \mathbb{E} \left[e^{\beta_{\text{SHE}} \int_{1-t}^1 \xi(s, B_s) - \frac{1}{2} \beta_{\text{SHE}}^2 t} \mid B_{1-t} = x \right]$$